# Nature-inspired Analog Computing on Silicon

Tetsuya ASAI and Yoshihito AMEMIYA Division of Electronics and Information Engineering Hokkaido University

#### Abstract

We propose CMOS analog circuits that emulate a model of the Belousov-Zhabotinsky (BZ) reaction, called the Oregonator. The BZ reaction gives us a lot of insight into developing new architectures based on a new paradigm in computing such as active wave computing. The proposed unit circuit exhibits both excitatory and oscillatory behaviors with very stiff responses, as observed in typical BZ reactions. We show spatiotemporal behaviors of an array of unit circuits, including synchronization of cell circuits and production of traveling waves, by using a simulation program of integrated circuit emphasis (SPICE).

### 1 Introduction

The formation of spatial and temporal patterns in dissipative and autocatalytic reaction systems have been spotlighted since almost every natural phenomenon, from traditional convective phenomena to modern neuro dynamics, can be categorized into these systems. Although a number of theoretical and numerical studies have been conducted to reveal the mechanism of those systems, the essential behaviors are still unknown due to their complexity and requirement of massive computational power in the numerical simulations. In this report, aiming at the development of high-speed emulators that advance understanding of such systems, we propose an analog CMOS circuit that implements a reaction-diffusion (RD) system [1].

Implementing RD systems in hardware (VLSI) has several merits. First, hardware RD system is very useful for simulating RD phenomena, even if the phenomena never occurs in nature. This implies that the hardware system is one possible candidate for developing an artificial RD system that is superior to natural system. Second, hardware RD system can operate much faster than actual RD systems. For instance, the velocity of chemical waves in Belousov-Zhabotinsky (BZ) reaction is  $O(10^{-2})$  m/s [2], while that of hardware RD system will be over a million times faster than that of the BZ reaction, independently of system size [3, 4, 5]. This property is useful for developers of RD applications because every RD application benefits from the operation speed.

### 2 The Reaction-Diffusion System

A reaction diffusion (RD) system is described by a set of partial differential equations

$$
\frac{\partial x_i(\mathbf{r},t)}{\partial t} = D_i \nabla^2 x_i(\mathbf{r},t) + f_i(x_i(\mathbf{r},t)),
$$
\n(1)



Figure 1: Nullclines and trajectories of the Oregonator in (a) oscillatory mode and (b) excitatory mode.

where **r** represents the space, t the time,  $\nabla^2$  the spatial Laplacian,  $D_i$  the diffusion constant,  $f_i$  the nonlinear reactive functions that depend on several different reactive species  $x_i$ . One well-known reactive function is described in a two-variable Oregonator, which is derived from the BZ reaction  $[1, 6]$ . The point dynamics are given by

$$
\frac{dx_1}{dt} = \frac{1}{\tau} \Big( x_1 \left( 1 - x_1 \right) - a \ x_2 \ \frac{x_1 - b}{b + x_1} \Big),\tag{2}
$$

$$
\frac{dx_2}{dt} = x_1 - x_2,\t\t(3)
$$

where  $x_1$  and  $x_2$  represent the concentration of HBrO<sub>2</sub> and Br<sup>−</sup> ions, respectively, while  $\tau$ , a and b represent the reaction parameters. The value of  $\tau$  is generally set at  $\tau \ll 1$ since the reaction rate of HBrO<sub>2</sub> ion is much faster than that of Br<sup>−</sup> ions. The nullclines of the Oregonator where  $dx_1/dt = 0$  and  $dx_2/dt = 0$  are given by

$$
x_2 = \frac{x_1 (x_1 + b)(1 - x_1)}{a (x_1 - b)}, \quad (\equiv l_1)
$$
 (4)

$$
x_2 = x_1. \quad (\equiv l_2) \tag{5}
$$

A cross point of those two nullclines  $(l_1 \text{ and } l_2)$  represents a fixed point of the Oregonator.

Figure 1 shows the nullclines and trajectories of the Oregonator with typical parametervalues ( $\tau = 10^{-2}$  and  $b = 0.02$ ). The value of parameter a was set at 1 [Fig. 1(a)] and 3 [Fig. 1(b)]. Depending on the position of the fixed point, the Oregonator exhibits oscillatory or excitatory behavior. When  $a = 1$ , the fixed point is located on nulllcline  $l_1$  at which  $dx_2/dx_1 > 0$ . In this case, the Oregonator exhibits limit-cycle oscillations [Fig. 1(a)]. The oscillation represents periodic oxidation-reduction phenomena of the BZ reaction. On the other hand, the fixed point is located on nulllcline  $l_1$  at which  $dx_2/dx_1 < 0$  when  $a = 3$ . Under this condition, the Oregonator exhibits excitatory behavior [Fig. 1(b)] and is stable at the fixed point as long as external stimulus is not given.



Figure 2: Nullclines and trajectories of the proposed analog cell in (a) oscillatory mode and (b) excitatory mode.

In the Oregonator, three circulative states are introduced according to the phase of oscillation; i.e., inactive (A), active  $(B \to C)$ , and refractory periods  $(D \to A)$ , as labelled in Fig. 1(b). The inactive, active, and refractory states represent a depletion in the Br<sup>−</sup> ion, an autocatalytic increase in the  $HBrO<sub>2</sub>$  ion (oxidation of the catalyzer), and a depletion in the Br<sup>−</sup> ion (reduction of the catalyzer), respectively. When the Oregonator is inactive, it easily become active  $(A \rightarrow B)$  by external stimuli. Then, it turns in refractory state  $(C \rightarrow D)$ . During the refractory state, the Oregonator can not be activated even if the external stimuli was given.

## 3 Analog CMOS Circuits for the BZ Reaction

We here propose a novel analog cell that is qualitatively equivalent to the Oregonator. We define the dynamics of a cell as

$$
\frac{dx_1}{dt} = \frac{1}{\tau} \Big( -x_1 + f(x_1 - x_2, \ \beta_1) \Big), \tag{6}
$$

$$
\frac{dx_2}{dt} = -x_2 + f(x_1 - \theta, \beta_2),
$$
\n(7)

where  $f(\cdot)$  represents a sigmoid function defined by

$$
f(x, \beta) = \frac{1 + \tanh \beta x}{2}.
$$
 (8)

The cell dynamics are designed so that the shape of nullclines and flows  $(\dot{x}_1, \dot{x}_2)$  are qualitatively equivalent to that of the Oregonator. The cubic nullcline  $(l_1$  in Fig. 1) is approximated by a nullcline of Eq. (6) as

$$
x_2 = x_1 - \beta_1^{-1} \tanh^{-1}(2x_1 - 1), \ (\equiv L_1)
$$
 (9)

while the linear nullcline  $(l_2 \text{ in Fig. 1})$  is approximated by a nullcline of Eq. (7) as

$$
x_2 = f(x_1 - \theta, \beta_2). \quad (\equiv L_2)
$$
\n
$$
(10)
$$



Figure 3: Numerical results of a RD system using the analog cell in (a) excitatory mode and (b) oscillatory mode.

An analog cell, whose dynamics are described by Eqs. (6) and (7), is very suitable for analog VLSI implementation because the sigmoid function can easily be implemented on the VLSIs by using differential-pair circuits.

The proposed cell exhibits qualitatively equivalent behaviors to the Oregonator, as shown in Fig. 2. The values of parameters are  $\tau^{-1} = 10$ ,  $\beta_1 = 5$  and  $\beta_2 = 10$ . When  $\theta = 0.5$ , the fixed point exists on a nulllcline [Eq. (9)] where  $dx_2/dx_1 > 0$ , and the system exhibits limit-cycle oscillations [Fig. 2(a)]. On the other hand, the system exhibits excitatory behavior [Fig. 2(b)] when the fixed point exists on a nulllcline (10) where  $dx_2/dx_1 < 0$  [Fig. 2(b)].

Now, let us introduce the cell dynamics into the RD model aiming at the constructing 2-D RD system. The dynamics of the RD system are described by substituting the reactive term in Eq.  $(1)$  with the right terms of Eqs.  $(6)$  and  $(7)$ . The discrete expression of the RD system is given by

$$
\frac{du_{i,j}}{dt} = \frac{1}{\tau} \Big( -u_{i,j} + f(u_{i,j} - v_{i,j}, \beta_1) \Big) + g_{i,j}^u, \tag{11}
$$

$$
\frac{dv_{i,j}}{dt} = -v_{i,j} + f(u_{i,j} - \theta, \ \beta_2) + g_{i,j}^v,\tag{12}
$$



Figure 4: An analog cell circuit consisting of single capacitor and two operationaltransconductance amplifiers (OTAs).



Figure 5: Layout of the analog cell designed with  $1.5-\mu$ m CMOS process (cell size: 70  $\times$  70  $\mu$ m<sup>2</sup>).

where the system variable  $x_i$  is replaced by  $u_{i,j}$  and  $v_{i,j}$ , while  $g_{i,j}^u$  and  $g_{i,j}^v$  represent external inputs to the cell (interactions between a cell and its neighboring cells) as

$$
g_{i,j}^u = D_u \frac{u_{i-1,j} + u_{i+1,j} + u_{i,j-1} + u_{i,j+1} - 4u_{i,j}}{h^2},
$$
  

$$
g_{i,j}^v = D_v \frac{v_{i-1,j} + v_{i+1,j} + v_{i,j-1} + v_{i,j+1} - 4v_{i,j}}{h^2}.
$$

Figure 3 shows spatiotemporal activities of the 2-D RD system with  $50 \times 50$  cells  $(\beta_1 = 5, \beta_2 = 10, h = 0.01 \text{ and } D_v = 0)$  where the values of  $v_{i,j}$  are represented in grayscale  $(v_{i,j} = 0$ : black,  $v_{i,j} = 1$ : white). The Neumann boundary condition was applied at the side of the square reaction-space. When  $\tau^{-1} = 10^2$  and  $\theta = 0.15$ at which the cell exhibits excitatory behavior, the 2D system produced target patterns [Fig. 3(a)], as observed in the basic RD system with the Oregonators. In the simulation, diffusion coefficient  $D_u$  was set at 3 ×10<sup>-4</sup>. The results indicate that the proposed RD system is qualitatively equivalent to the basic RD system with the Oregonators since the excitatory property of the analog cells is inherently the same as that of the Oregonator.



Figure 6: Experimental results of the BZ cell

Figure 3(b) shows dynamic behaviors of the 2-D RD system with  $D_u = 10^{-4}$  and  $\theta = 0.5$  at which the cell exhibits oscillatory behavior. Initial values of the cells were randomly chosen as  $u_{i,j} = \text{RAND}[0,1]$  and  $v_{i,j} = \text{RAND}[0,1]$ . The system produced 2-D phase-lagged stable synchronous patterns called modelock, due to the weak coupling between the neighboring cells. When  $D_u > 10^{-3}$ , all cells exhibit synchronous oscillation, namely, no spatial pattern was produced.

An analog circuit of the proposed cell and its device layout are shown in Figs. 4 and 5, respectively. The circuit consists of single capacitor and two operationaltransconductance amplifiers (OTAs), which implies that the circuit can easily be implemented on silicon VLSIs using conventional CMOS technology. The circuit can be obtained by qualitative approximation of Eqs. (6) and (7).

When the rate constant of Eq. (6) is much larger than that of Eq. (7), the differential term of Eq. (6) can be neglected ( $\tau \ll 1$ ), as introduced in Sec. II. On the other hand, Eq. (7) with  $\beta_2 \to \infty$  forces the values of variable  $x_2$  to be 0 when  $x_1 \leq \theta$ , while  $x_2 \to 1$ when  $x_1 > \theta$ . If the variable  $x_2$  is forced to have the value within [0:1], the temporal difference in Eq. (7) can approximately be represented by binary values. Consequently, we obtain a new dynamical equation from Eqs. (6) and (7) as

$$
x_1 = f(x_1 - x_2, \beta_1) \tag{13}
$$

$$
\frac{dx_2}{dt} = \begin{cases} w & (\text{if } x_1 > \theta) \\ -w & (\text{else}) \end{cases}
$$
(14)

where w represent positive and small constant. In Fig. 4, an OTA labeled as  $\beta_1$  serves as the function of Eq.  $(13)$ , while a capacitor C and the rest OTA receiving voltage θ produce the dynamics for Eq. (14). The positive constant w is implemented in the OTA (labeled as  $w$ ) where  $w$  corresponds to the source current of a differential pair. The output current of the OTA  $(w)$  becomes zero when the voltage of output node  $x_2$ is equal to the supply voltage (VDD or VSS). The value of  $x_2$  is thus restricted within [VDD:VSS].



Figure 7: Chip photograph of 2-D BZ circuits.

Figure 6 shows experimental results of the fabricated chip that implements the reaction circuit shown in Fig. 5 (MOSIS 1.5- $\mu$ m CMOS, cell size:  $70 \times 70 \ \mu$ m<sup>2</sup>). Supply voltages of an OTA of  $\beta_1$  was set at VDD = 4 V and VSS = 0.5 V, while that of the rest OTA (w) was set at VDD = 5 V and VSS = GND. The threshold  $\theta$  was set at 2.5 V so that the circuit exhibits oscillatory behaviors. In the device layout shown in Fig. 5, the capacitor  $C$  was implemented by a MOS capacitor (lower-right rectangle in Fig. 5). As expected, the circuit exhibited qualitatively same behaviors as the Oregonator; i.e., stiff nonlinear oscillations.

#### 4 Summary

We proposed an analog CMOS circuit that that imitates a model of Belousov-Zhabotinsky (BZ) reaction. Numerical simulations and experimental results of the fabricated chip showed that the circuit can successfully produce excitatory responses and thus spiral waves in the same way as natural reaction-diffusion (RD) systems. These results encourage us to develop novel applications based on natural RD phenomena using the hardware RD devices.

The proposed devices and circuits are useful not only for the hardware RD system but also for constructing modern neuro-chips. The excitatory and oscillatory behaviors of the RD circuit are very similar to actual neurons that produce sequences in time of identically shaped pulses, called spikes. Recently, Fukai showed that an inhibitory network of spiking neurons achieves robust and efficient neural competition on the basis of a novel timing mechanism of neural activity [7]. A network with such a timing mechanism may provide an appropriate platform for the development of analog VLSI circuits that overcome the problems of analog devices, namely their lack of precision and reproducibility.

# References

- [1] G. Nicolis and I. Prigogine, Self-organization in Nonequilibrium Systems From Dissipative Structures to Order through Fluctuations. John Wiley & Sons, Inc., 1977.
- [2]  $\AA$ . Tóth, V. Gáspár, and K. Showalter: J. Phys. Chem **98** (1994) 522.
- [3] T. Asai, Y. Nishimiya, and Y. Amemiya: IEICE Trans. Fundamentals E85-A (2002) 2093.
- [4] T. Asai, Y. Nishimiya, and Y.Amemiya: Proceedings. Int. Semiconductor Device Research Symp (2001), 141-4.
- [5] Y. Nishimiya, T. Asai, and Y.Amemiya: Ext. Abst. 2001 Int. Conf. on Solid State Devices and Materials (2001), 404-5.
- [6] R. J. Field and M. Burger, Oscillations and travelling waves in chemical systems. John Wiley & Sons, Inc., 1985.
- [7] T. Fukai: Biol. Cybern. 75 (1996) 453.